

You Can Always Get What You Want: The Impact of Prior Assumptions on Interpreting GW190412

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ABSTRACT

GW190412 is the first observation of a black hole binary with definitively unequal masses. GW190412’s mass asymmetry, along with the measured positive effective inspiral spin, allowed for inference of a component black hole spin: the primary black hole in the system was found to have a dimensionless spin magnitude between 0.17 and 0.59 (90% credible range). We investigate how the choice of priors for the spin magnitudes and tilts of the component black holes affect the robustness of parameter estimates for GW190412, and report Bayes factors across a suite of prior assumptions. Depending on the waveform family used to describe the signal, we find either marginal to moderate (2:1–7:1) or strong ($\gtrsim 20:1$) support for the primary black hole being spinning compared to cases where only the secondary is allowed to have spin. We show how these choices influence parameter estimates, and find the asymmetric masses and positive effective inspiral spin of GW190412 to be qualitatively, but not quantitatively, robust to prior assumptions. Our results highlight the importance of both considering astrophysically-motivated or population-based priors in interpreting observations, and considering their relative support from the data.

1. INTRODUCTION

GW190412 (Abbott et al. 2020) was the first reported observation of a binary black hole (BBH) from the third observing run (O3) of the Advanced LIGO (Aasi et al. 2015) and Advanced Virgo (Acernese et al. 2015) detector network. GW190412’s source is the first system to have definitively unequal masses (cf. Abbott et al. 2019a), with the primary black hole (BH) being $\sim 30 M_\odot$ and the secondary BH being $\sim 8 M_\odot$. In addition to unveiling emission from higher-order multipoles (HMs), this asymmetry allowed for enhanced constraints on the intrinsic and extrinsic parameters of the BBH system.

The spins of compact binary components are difficult to measure from gravitational-wave (GW) signals (Poisson & Will 1995; Vitale et al. 2014; Pürrer et al. 2016; Abbott et al. 2016a). Typically, spin constraints are presented in terms of mass-weighted combinations of the two component spins: the effective inspiral spin

$$\chi_{\text{eff}} = \frac{m_1 \chi_1 \cos \theta_1 + m_2 \chi_2 \cos \theta_2}{m_1 + m_2}, \quad (1)$$

where $m_1 \geq m_2$ are the component masses, χ_i are the dimensionless spin magnitudes, and θ_i are the angles between the spins and the Newtonian orbital angular momentum \vec{L} , encodes information about the spin components aligned with the orbital angular momentum (Damour 2001; Racine 2008; Santamaría et al. 2010; Ajith et al. 2011), whereas in-plane spins are characterized by the effective precession spin (Schmidt et al. 2015)

$$\chi_p = \max \left\{ \chi_1 \sin \theta_1, \frac{q(4q+3)}{4+3q} \chi_2 \sin \theta_2 \right\}. \quad (2)$$

The LIGO Scientific & Virgo Collaboration (LVC) reported an effective spin for GW190412 of $\chi_{\text{eff}} = 0.25^{+0.08}_{-0.11}$ (median and 90% credible interval; Abbott et al. 2020). Since χ_{eff} is positive and constrained away from zero, at least one of the BHs in the GW190412 system had a spin direction in the same hemisphere as \vec{L} during the GW inspiral. GW190412 also exhibited marginal hints of orbital precession, which is consistent with at least one of the BH spins being non-zero.

A BBH with $\chi_{\text{eff}} > 0$ has been observed before in GW151226 (Abbott et al. 2016b; Miller et al. 2020), and potentially in GW170229 (Abbott et al. 2019a; Chatzi-

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ioannou et al. 2019). However, the larger mass of the primary BH in GW190412 relative to the secondary BH allowed for the spin of the primary to be inferred as $\chi_1 = 0.43^{+0.16}_{-0.26}$. This is because when $m_1 \gg m_2$ the primary spin is much more important in determining the dynamics of the system (as illustrated by the mass weighting in χ_{eff} and χ_p), and we are less sensitive to the value of the secondary spin. GW190412 therefore is the first high-significance detection of a compact binary system with an observable component spin.¹

GW190412’s high primary spin may be difficult to reconcile with theoretical modeling of massive binary stars in isolation. Detailed modeling of core–envelope interaction in massive stars finds angular momentum transport to be highly efficient, driving stellar cores to extremely slow rotation prior to their collapse into a compact object (Kushnir et al. 2016; Zaldarriaga et al. 2018; Fuller et al. 2019; Fuller & Ma 2019; Belczynski et al. 2020). This theoretical underpinning is corroborated by the majority of LVC observations, which are mostly consistent with $\chi_{\text{eff}} \approx 0$ (Abbott et al. 2019b; Safarzadeh et al. 2020; Miller et al. 2020). Though the birth spins of some BHs in high-mass X-ray binaries have been interpreted as near extremal ($\chi \approx 1$; see Miller & Miller 2015 and references therein), it is unclear whether these systems will ever evolve to be BBHs that merge within a Hubble time (e.g., Belczynski et al. 2012; Qin et al. 2019). Following this reasoning, multiple groups have proposed that the high spin of the primary BH in GW190412 is the result of an alternative formation scenario to canonical isolated binary evolution, such as dynamical assembly in young star clusters (Di Carlo et al. 2020), hierarchical mergers in massive stellar clusters (Kimball et al. 2020; Gerosa et al. 2020; Rodriguez et al. 2020), active galactic nucleus (AGN) discs (Tagawa et al. 2020), Population III stars (Kinugawa et al. 2020), and mergers induced from the secular evolution in hierarchical systems (Hamers & Safarzadeh 2020).

On the other hand, the second-born BH in BBH merger progenitors can be significantly spun up through tidal locking of the stellar core with the first-born BH (Qin et al. 2018; Bavera et al. 2020). If GW190412 could instead be explained by a highly spinning secondary BH, the standard isolated formation scenario with a low-spinning primary could again be viable. To

this end, Mandel & Fragos (2020) provide a reinterpretation of the LVC analysis using a prior motivated by theoretical predictions of BBH progenitors formed in isolation. Assuming a prior with a zero-spin primary BH and a secondary BH whose spin projection is aligned with the orbital angular momentum, Mandel & Fragos (2020) reweight the public posterior samples of GW190412 (LIGO Scientific Collaboration and Virgo Collaboration 2020), effectively interpreting the measured value of χ_{eff} as originating from the secondary’s spin rather than the primary’s. To compensate for the non-zero effective spin of GW190412, the reweighted posteriors from this analysis point to a highly-spinning secondary BH with $\chi_2 \gtrsim 0.64$ (Mandel & Fragos 2020).

Though predictions for the formation rate of these systems are highly sensitive to the uncertain prescription for natal BH spins, recent work has found that for systems with asymmetric masses such as GW190412, the highly-spinning secondary BH interpretation is more probable from an isolated evolution standpoint than a moderately-spinning primary (e.g., Olejak et al. 2020). This is consistent with the current catalog of GWs, since individual spins are poorly constrained in all previously observed BBHs (Abbott et al. 2019b). However, even this formation mechanism struggles to accommodate GW190412, as systems where the secondary BH has been significantly spun up due to tidal interactions have short merger timescales and a merger rate in the local universe that is at least an order of magnitude lower than what is estimated for GW190412-like systems (Safarzadeh & Hotokezaka 2020).

Nonetheless, while various assumptions may be made to represent the prior belief for parameters given an astrophysical model, it is critical to determine whether a given model is supported by the data. The amount by which the data supports a specific model (in this work, a prior) is encoded in the Bayesian evidence. While varying prior assumptions will yield differing parameter estimates, the ratio of evidences between models—the *Bayes factor* \mathcal{B} —indicates whether any one prior assumption is favored or disfavored by the data compared to another. This is particularly important to verify for the case of strong priors, since they might drive the posteriors to potentially arbitrary values at the expense of the evidence: *if you torture the data long enough, it will confess to anything* (Coase 1982). For example, in the analysis of GW151226 (Abbott et al. 2016b), Vitale et al. (2017a) showed how if one uses a prior that enforces small (~ 0.1) spin magnitudes, the evidence decreases by a factor of 50 compared to a uniform prior, while the posteriors still look reasonable. It is only by

¹ A potential BBH with a highly spinning primary component was reported in Zackay et al. (2019), though the astrophysical origin of the signal is under debate (Nitz et al. 2019; Ashton & Thrane 2020), and due to the low signal-to-noise of this event, the spin interpretation depends heavily on the choice of prior (Nitz et al. 2019; Galadage et al. 2019; Huang et al. 2020).

comparing evidences, i.e. calculating Bayes factors, that one can assess if a model matches the data.

In this paper, we explore various prior assumptions for the interpretation of GW190412 and calculate Bayes factors between these model assumptions. The priors we choose are motivated by various astrophysical models presented in the literature, with a particular focus on the spin of the second-born BH, and the astrophysically relevant question of whether the primary is spinning.

In Sec. 2 we explain the various prior assumptions we choose when analyzing the data, and their astrophysical motivation. We present Bayes factors across these prior assumptions in Sec. 3, and examine the impact of differing prior assumptions on the parameter estimation for GW190412 in Sec. 4. In Sec. 5 we discuss the results of our analysis and their impact on the interpretation of GW190412, and comment on astrophysical implications.

2. DATA ANALYSIS AND PRIOR ASSUMPTIONS

To investigate the impact of prior assumptions on the inferred parameters of GW190412 and the Bayes factors between these assumptions, we perform parameter estimation using a suite of prior assumptions motivated by various astrophysical predictions. We use the publicly-available data for GW190412 (LIGO Scientific Collaboration and Virgo Collaboration 2020) and follow the parameter-estimation procedure used in Abbott et al. (2020). Our results are produced using a highly parallelized version of BILBY (Ashton et al. 2019; Smith et al. 2019; Romero-Shaw et al. 2020), which computes posterior probability distributions for the properties of the source as well as model evidences.

We use both the *Phenom* and *EOB* families of waveform approximants in our analysis. We use *IMRPhenomPv3HM* (Khan et al. 2019, 2020) and *SEOBNRv4PHM* (Pan et al. 2014; Babak et al. 2017; Ossokine et al. 2020), both of which include the effects of spin precession and HM moments. Inclusion of HMs in waveform approximants is crucial for the parameter estimation of GW190412, as this more complete physical picture of the GW signal is necessary to accurately constrain the mass ratio ($q = m_2/m_1 < 1$) and spins (Van Den Broeck & Sengupta 2007; Graff et al. 2015; Calderón Bustillo et al. 2016; Varma & Ajith 2017; Abbott et al. 2020). Systematic differences are expected between analyses using *Phenom* and *EOB* approximants, as evident in Abbott et al. (2020). Though we use the *Phenom* approximant for all 7 prior configurations described below, due to the computational cost of the *EOB* approximant we only run this with two exemplary prior configurations.

The priors we consider are:

- A) Uniform in spin magnitude for both components, isotropic and unconstrained in spin tilts. This uninformative prior is used in Abbott et al. (2020); it does not make strong assumptions about spin orientations or magnitudes, and its broad support enables reweighting by different priors (e.g., Thrane & Talbot 2019; Mandel et al. 2019).
- B) Uniform in spin magnitude and isotropic in spin tilt for the primary BH, with a non-spinning secondary. A spinning primary and a non-spinning secondary may be expected if BHs are born with small spins, but the larger BH is the result of a previous BH merger and has gone on to form a new binary in a dense stellar environment such as a globular or nuclear cluster (Fishbach et al. 2017; Gerosa & Berti 2017; Rodriguez et al. 2019; Kimball et al. 2020; Gerosa et al. 2020). In this scenario, we would typically expect the primary spin magnitude to be $\chi_1 \sim 0.67$.
- C) Non-spinning primary BH with unconstrained spin for the secondary BH. This is representative of an isolated formation scenario, with a secondary that can be spun up through tidal interactions (Qin et al. 2019; Bavera et al. 2020). The unconstrained spin tilt, however, allows for significantly misaligned spins, which are difficult to attain for BBHs in the standard isolated evolution scenario (e.g., Kalogera 2000; Fryer et al. 2012; Rodriguez et al. 2016).
- D) Same as Model C, but with spin tilts constrained to be in the same hemisphere as the orbital angular momentum: $\vec{\chi}_2 \cdot \vec{L} \geq 0$ ($0^\circ \leq \theta_2 \leq 90^\circ$). This is similar to the prior assumption used in Mandel & Fragos (2020).
- E) Same as Model C, but with spin tilts for the secondary constrained to $0^\circ \geq \theta_2 \geq 10^\circ$. This model has been used to represent near-aligned spins (e.g., Vitale et al. 2017b), as predicted from the coevolution of isolated binaries and weak BH natal kicks at birth.
- F) Same as Model C, but with spin tilts for the secondary perfectly aligned with the orbital angular momentum ($\theta_2 = 0^\circ$).
- G) Non-spinning primary and secondary. This is an extreme assumption that we expect will struggle to match the data due to the positive measured χ_{eff} and marginal precessional information.

These configurations are summarized on the left side of Table 1. For all other parameters, we use priors analogous to those used by the LVC in the analysis of GW190412 (Abbott et al. 2020).

3. BAYES FACTORS

Given the observation of GW190412, we can identify which astrophysical model is best supported by the data. To quantify the relative support for different models, we would ideally use the *odds ratio*; the odds ratio between models \mathcal{M}_i and \mathcal{M}_j is defined as

$$\mathcal{O}_{\mathcal{M}_i, \mathcal{M}_j} = \frac{p(\mathcal{M}_i|d)}{p(\mathcal{M}_j|d)}, \quad (3)$$

where $p(\mathcal{M}_i|d)$ is the posterior probability of model \mathcal{M}_i given the data d . Using Bayes' theorem, the odds ratio can be written as

$$\mathcal{O}_{\mathcal{M}_i, \mathcal{M}_j} = \frac{p(\mathcal{M}_i)}{p(\mathcal{M}_j)} \frac{p(d|\mathcal{M}_i)}{p(d|\mathcal{M}_j)}. \quad (4)$$

Here the first term is the prior odds: our expectation for the relative probabilities of the two models before observing the data. For example, predictions for the local BBH merger rate from isolated binary evolution range from $\sim 8\text{--}200 \text{ Gpc}^{-3} \text{ yr}^{-1}$ (e.g., Eldridge et al. 2017; Klencki et al. 2018; Giacobbo & Mapelli 2018, 2020) while predictions for the local BBH merger rate through dynamical assembly in globular clusters ranges from $\sim 0.8\text{--}35 \text{ Gpc}^{-3} \text{ yr}^{-1}$ (e.g., Fragione & Kocsis 2018; Hong et al. 2018; Rodriguez & Loeb 2018); thus, from the ratio of these predicted rates one may estimate a prior odds between the two channels of $\sim 0.2\text{--}250$. In addition to considering expected rates, prior odds could also factor in our belief in the accuracy of different physical prescriptions, for example the efficiency of angular momentum transport in massive stars. Given the uncertainties in the prior odds, we concentrate on the second term in Eq. (4), the Bayes factor: the ratio of evidences for the two models.

The model evidence, or marginalised likelihood, is

$$p(d|\mathcal{M}_i) = \int d\boldsymbol{\vartheta}_i \mathcal{L}(\boldsymbol{\vartheta}_i) p(\boldsymbol{\vartheta}_i|\mathcal{M}_i), \quad (5)$$

where the integral is over the parameters $\boldsymbol{\vartheta}_i$ describing our source (masses, spins, etc.), $\mathcal{L}(\boldsymbol{\vartheta}_i) = p(d|\boldsymbol{\vartheta}_i)$ is the likelihood of the parameters (Cutler & Flanagan 1994), and $p(\boldsymbol{\vartheta}_i|\mathcal{M}_i)$ is our prior probability density on the parameters within model \mathcal{M}_i , as described in Sec. 2. Thus, the Bayes factor is given by

$$\mathcal{B}_{\mathcal{M}_i, \mathcal{M}_j} = \frac{\int d\boldsymbol{\vartheta}_i p(d|\boldsymbol{\vartheta}_i) p(\boldsymbol{\vartheta}_i|\mathcal{M}_i)}{\int d\boldsymbol{\vartheta}_j p(d|\boldsymbol{\vartheta}_j) p(\boldsymbol{\vartheta}_j|\mathcal{M}_j)}. \quad (6)$$

When considering models with more parameters, or with parameters allowed to vary on a larger domain, we expect that we may be able to fit the data better, giving higher likelihoods. In calculating evidences, this is counterbalanced by the increased prior volume: as we spread the total prior probability (which must integrate to 1) over a larger volume where the likelihood can have potentially negligible support, its density around the maximum likelihood region may decrease, resulting in a lower evidence. This *Occam factor* allows the Bayes factor to be used to determine if more complicated models are needed to explain data (MacKay 2003, Chapter 28).

When considering spins measured with GW observations, we are typically only sensitive to particular mass-weighted combinations of the 6 spin degrees of freedom (cf. Poisson & Will 1995; Chatziioannou et al. 2014; Pürrer et al. 2016; Vitale et al. 2017a). Therefore, it may be possible to fit the data well by assuming only a single component is spinning, and we would not anticipate a strong preference in favor of a more complicated model including two spinning bodies. In cases when there is a large asymmetry in masses the secondary spin may become irrelevant, and the properties of the signal may be completely determined by the primary spin. When the secondary spin has negligible impact on the likelihood, we expect there will be no preference between models with and without a secondary spin as it is unconstrained and its introduction incurs no Occam factor penalty.²

In Table 1 we show Bayes factors for each prior compared to the standard LVC prior (Model A) using the **Phenom** approximant and selected results with the **EOB** approximant. Model A is preferred over the other prior models; the extra freedom allowed by having two spinning bodies enables a better fit to the data (as illustrated by the maximum likelihood value), and this improvement is sufficient to overcome the Occam factor from the larger prior volume.

Despite the significant asymmetry in masses, the secondary spin still has some impact on the signal, as can be seen by comparing the Bayes factor between Model A and Model B (1.6:1 with the **Phenom** approximant).

We find marginal to moderate support for Model A relative to prior configurations where only the secondary is spinning (Models C–F) with the **Phenom** approximant, and strong support for Model A relative to non-spinning

² Analogously, when spin precession is not measurable, such that the posterior distribution for χ_p is identical to the prior, we expect no preference between using a waveform approximant that includes spin precession and one that only includes the effects of the spin components aligned with the orbital angular momentum, assuming that the priors on the aligned components of the spin are equivalent.

Table 1. *Left:* Prior assumptions for component spin magnitudes and spin tilts in each model. We use short-hand for the distributions we consider: δ uses a fixed value, U denotes a uniform distribution, and ISO an isotropic distribution (uniform in $\cos(\theta)$). Angular assumptions are omitted when the spin magnitudes for that component are forced to zero. *Right:* The maximum value of the log likelihood ($\log_{10} \mathcal{L}_{\max}$) and Bayes factors (\mathcal{B}) for each model. Each Bayes factor is calculated relative to the uninformative LVC prior for the respective waveform approximant (Model A and Model A-EOB for **Phenom** and **EOB**, respectively). For reference, $\mathcal{B} < 1$ ($\log_{10} \mathcal{B} < 0$) means that data prefers the reference model, $\mathcal{B} \gtrsim 3:1$ ($\log_{10} \mathcal{B} \gtrsim 0.5$) indicates moderate evidence for the new hypothesis, and $\mathcal{B} \gtrsim 10:1$ ($\log_{10} \mathcal{B} \gtrsim 1.0$) indicates strong evidence for the new hypothesis. The rightmost column gives the estimated 1-sigma uncertainty in $\log_{10} \mathcal{B}$.

Model	Prior assumption				Evidence			
	χ_1	θ_1	χ_2	θ_2	$\log_{10} \mathcal{L}_{\max}$	\mathcal{B}	$\log_{10}(\mathcal{B})$	$\sigma_{\log_{10}(\mathcal{B})}$
A	U[0, 0.99]	ISO[0°, 180°]	U[0, 0.99]	ISO[0°, 180°]	77.0	1.0	0.00	0.08
A-EOB	U[0, 0.99]	ISO[0°, 180°]	U[0, 0.99]	ISO[0°, 180°]	77.1	1.0	0.00	0.10
B	U[0, 0.99]	ISO[0°, 180°]	$\delta(0)$	—	76.5	6.2×10^{-1}	-0.20	0.09
C	$\delta(0)$	—	U[0, 0.99]	ISO[0°, 180°]	75.6	1.5×10^{-1}	-0.80	0.09
D	$\delta(0)$	—	U[0, 0.99]	ISO[0°, 90°]	75.5	3.6×10^{-1}	-0.44	0.09
D-EOB	$\delta(0)$	—	U[0, 0.99]	ISO[0°, 90°]	74.4	4.9×10^{-2}	-1.30	0.10
E	$\delta(0)$	—	U[0, 0.99]	ISO[0°, 10°]	75.3	5.3×10^{-1}	-0.27	0.09
F	$\delta(0)$	—	U[0, 0.99]	$\delta(0^\circ)$	75.2	5.0×10^{-1}	-0.30	0.09
G	$\delta(0)$	—	$\delta(0)$	—	71.5	2.8×10^{-3}	-2.54	0.08

primary configurations when using the **EOB** approximant. For the non-spinning primary configurations and the **Phenom** approximant, we find the greatest support for Model E, which is only disfavored relative to our fiducial model by a Bayes factor of $\simeq 2:1$. As the opening angle for θ_2 increases, we see a decreasing trend in the Bayes factor that is likely due to the Occam factor suffered by the models with larger possible misalignment, since the maximum likelihood for these three models is relatively constant. With a non-spinning primary, the secondary BH needs to have significant spin aligned with the orbital angular momentum in order to match the observed signal. Therefore, the tilt is constrained to be small, and there is little in-plane spin. Though precession is possible in Models C–E, it is not possible to have a large χ_p given both the mass ratio and the need to match the χ_{eff} measurement. With the **Phenom** approximant, the case where we can draw the most confident conclusion is in comparison to the prior configuration with zero spins, which is disfavored by a Bayes factor of $> 400:1$.

We find strong support against the non-spinning primary hypothesis when using the **EOB** approximant. The maximum likelihood value using the LVC prior is greater than that of the non-spinning primary, aligned-spin secondary prior used in [Mandel & Fragos \(2020\)](#) by a factor of $\gtrsim 500$. Though the LVC prior configuration has a larger prior volume, the strong support in the data for a spinning primary leads to a Bayes factor of $\gtrsim 20:1$

relative to the non-spinning primary hypothesis. The Bayes factor between the **Phenom** and **EOB** approximants using the LVC prior (Models A and A-EOB) is $0.97:1$, indicating no preference for one of these approximants over the other. We discuss implications of these Bayes factors further in Sec. 5.

4. PARAMETER ESTIMATION

Prior assumptions can have a strong effect on the measurement of intrinsic and extrinsic parameters inherent to a BBH coalescence. Here, we investigate the robustness of parameter estimates for GW190412 across our various prior assumptions, with a particular focus on spin parameters.

4.1. Mass Ratio

GW190412 is the first BBH with definitely unequal masses, with a reported mass ratio at the 90% credible level of $0.25 \leq q \leq 0.45$ using the **Phenom** approximant and $0.21 \leq q \leq 0.31$ using the **EOB** approximant ([Abbott et al. 2020](#)). In Fig. 1 we show the posterior distributions for q across our different priors and waveform approximants. Aside from the (strongly disfavored) zero spin Model G, we find the mass ratio to be constrained to $q \lesssim 0.57$ at the 99% credible level.

There is a noticeable difference in the posterior distribution for q when using priors where the primary is spinning compared to those where only the secondary is spinning. We find that the posterior for q pushes to larger values when $\chi_1 = 0$, with a median of 0.39 (0.39)

in Model D compared to 0.32 (0.26) in Model A when using the **Phenom** (EOB) approximant. This change in q results in a more massive secondary that can more easily account for the observed effective spins.

4.2. Aligned Spin and Precession

When the primary is forced to be non-spinning the effective inspiral spin migrates to lower values (lower left panel of Fig. 1); with a non-spinning primary, χ_{eff} must always be less than q . Using the **Phenom** (EOB) approximant, we find χ_{eff} to be 0.06–0.22 (0.10–0.23) for the non-spinning primary Model D compared to 0.12–0.30 (0.19–0.36) for Model A at the 90% credible level. Using the source parameters derived with the LVC prior, the *largest* χ_{eff} that can be attained from a system with a non-spinning primary of mass $m_1 \sim 30 M_\odot$ and a spinning secondary with mass $m_2 \sim 8 M_\odot$ is $\chi_{\text{eff}}|_{\chi_1=0} \lesssim 0.21$. Thus, prior models with a non-spinning primary need to compensate by jointly increasing q and decreasing χ_{eff} . However, for all our models where at least one BH is spinning we find $\chi_{\text{eff}} > 0.08$ at the 90% credible level.

Considering in-plane spins, χ_p shows a larger variation between the prior models. This is expected, since χ_p affects the likelihood only mildly and our prior models put restrictions on spin tilts. For both waveform approximants, when the primary is non-spinning χ_p is $\lesssim 0.28$ at the 90% level, and rails against the physical boundary of $\chi_p = 0$ (consistent with no precession). The median posterior value for χ_p drops even more precipitously when a large degree of misalignment is not allowed; for Model E we recover a median χ_p of 0.027. Thus, if indeed the primary BH is non-spinning, the marginal hints of precession in GW190412 disappear and the system is consistent with having a perfectly aligned secondary spin.

4.3. Component Spins

In Fig. 2 we show marginalized posterior distributions for the two component spins, χ_1 and χ_2 . In the prior configurations where χ_1 is non-zero, we recover similar posterior distributions across the **Phenom** results, though when χ_2 is forced to zero the distribution shifts to slightly higher values with a median χ_1 that is 0.03 larger than in the LVC prior case. This is because the primary BH must now account for all spin effects in the data without a contribution from the secondary. The χ_1 posteriors are also consistent with the Bayes factors reported in Table 1 in favor of models where the primary BH is spinning: for Models A and C (which are nested since Model C can be obtained by fixing $\chi_1 = 0$ in Model A), the Bayes factor can be obtained by comparing the prior to the posterior at $\chi_1 = 0$ (Chatziioannou et al.

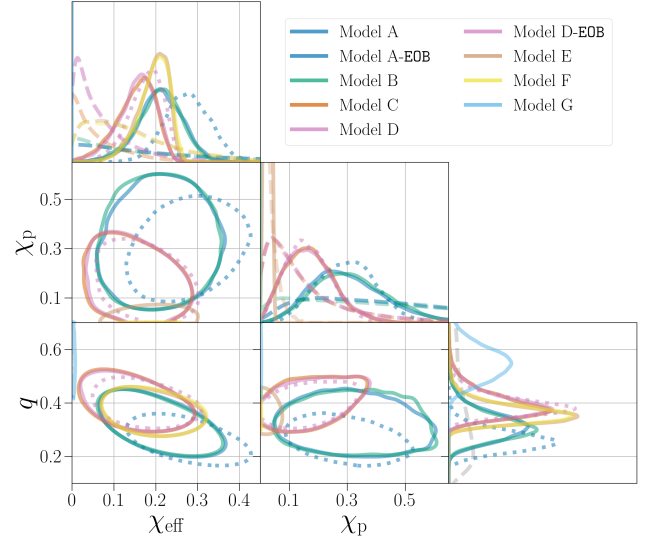


Figure 1. Joint and marginalized posterior distributions for the effective inspiral spin χ_{eff} , effective precession spin χ_p , and mass ratio q . The posteriors recovered for each prior configuration are shown using different colors. Posteriors using the **Phenom** approximant are shown in solid lines, and the subset of posteriors using the EOB approximant are shown in dotted lines. In the panels showing marginalized posterior distributions, the prior distributions for each configuration are shown with corresponding dashed lines; for q the prior distribution is the same for all configurations and we display it with a single gray dashed line. We see mild, yet noticeable differences in the posterior distributions for q and χ_{eff} when we constrain the primary spin to $\chi_1 = 0$, though we still recover asymmetric masses and a non-zero effective spin at high confidence for all runs with reasonable Bayes factors. For the non-spinning configuration (Model G) only $\chi_{\text{eff}} = 0$ and $\chi_p = 0$ are allowed.

2014). The top panel of Fig. 2 shows that the prior at $\chi_1 = 0$ is larger than the posterior for both waveform approximants, pointing to a Bayes factor in favor of a spinning primary.

We see larger variation in χ_2 across the prior configurations. The standard LVC prior recovers a broad, uninformative distribution in χ_2 . However, when χ_1 is forced to zero spin, χ_2 is constrained away from zero in all cases; in these models, we find χ_2 to be consistent with maximally spinning and have $\chi_2 \gtrsim 0.50$ at the 90% credible level (cf. Mandel & Fragos 2020). The EOB results push to slightly higher secondary spins than the **Phenom** results with non-spinning primary configurations, with $\chi_2 \gtrsim 0.62$ at the 90% credible level for Model D-EOB. In all cases where spin-tilt misalignment is allowed, we find a preference for some degree of misalignment in the spins; at the 90% credible level, we find $\theta_2 > 18.5^\circ$ in Model C, $\theta_2 > 18.1^\circ$ ($\theta_2 > 16.3^\circ$) in Model D (Model D-EOB), and $\theta_2 > 3.2^\circ$ in Model E. Thus, we

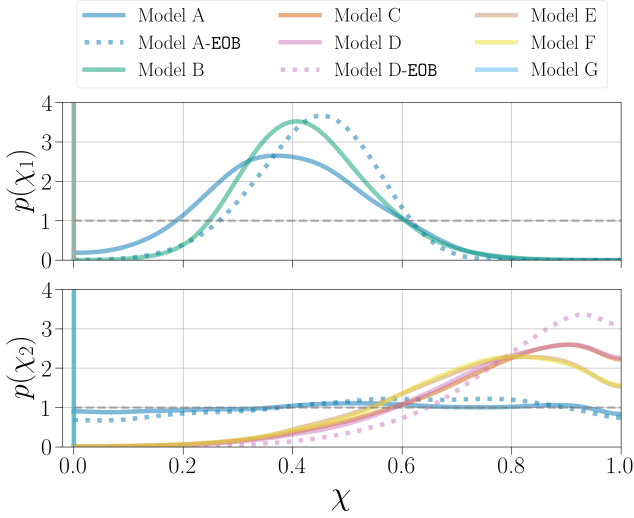


Figure 2. Marginalized spin distributions for the primary dimensionless BH spin (χ_1 ; *top*) and the secondary dimensionless BH spin (χ_2 ; *bottom*). Posteriors recovered for each prior configuration are in different colors, and the flat priors for each spin magnitude are shown as a gray dashed line. Posteriors attained using the **Phenom** approximant are shown in solid lines, and the subset of posteriors from the **EOB** approximant are shown in dotted lines. When the primary BH is non-spinning, χ_2 is constrained to higher values and consistent with maximally spinning ($\chi_2 = 1$). For priors where the primary or secondary BH are forced to be non-spinning, the posterior is shown by a vertical line at $\chi = 0$.

find that precession (albeit possibly immeasurable) is permitted in all prior models that allow for spin tilts.

5. DISCUSSION AND CONCLUSIONS

GW190412 is an astrophysically compelling event that resides in a previously unobserved region of BBH parameter space. The effective inspiral spin of the system indicates that at least one of the component BHs is spinning. This work investigates whether a spinning primary BH or a spinning secondary BH is better supported by the data, and how these hypotheses affect the inferred parameters of GW190412.

Our main results are summarized in Table 1. The broad LVC prior (Model A) with both BHs spinning is preferred over the other models, despite the larger prior volume. The degree of preference depends on the waveform approximant used, as the effect of waveform systematics are non-negligible for this event (Abbott et al. 2020). We recover marginal support in favor of Model A compared to the model where only the primary is spinning (Model B). When using the **Phenom** approximant we find marginal to moderate evidence in favor of Model A compared to models where only the secondary is allowed to spin (Models C–F), whereas with the **EOB**

approximant we find strong evidence in support of the LVC prior model compared to priors with a non-spinning primary. The data strongly supports Model A over the hypothesis where neither BH is spinning (Model G) for both waveform approximants.

The **Phenom** approximant gives broader parameter constraints than the **EOB** approximant in both Abbott et al. (2020) and this work. In Fig. 1, we see that the non-spinning primary prior models move the posterior distributions for q (χ_{eff}) to higher (lower) values to better allow the secondary to account for the spin information in the signal. This comes at the cost of matching the data, as the maximum likelihood values are $\lesssim 1.5$ dex lower for models where only the secondary is spinning compared to Model A. Whereas the **Phenom** approximant measures q (χ_{eff}) to be $\lesssim 0.45$ ($\lesssim 0.29$) and the 90% credible level, the **EOB** approximant recovers $\lesssim 0.31$ ($\lesssim 0.34$). The lower mass ratio and higher effective spin from the **EOB** analysis makes it more difficult for the data to accommodate a non-spinning primary, with maximum likelihood values that are $\simeq 2.7$ dex lower for the model where only the secondary is spinning. Despite the larger prior volume, we find a spinning primary hypothesis to be favored over a non-spinning primary hypothesis by a Bayes factor of $\gtrsim 20:1$.

The prospect of a non-spinning primary BH was explored in Mandel & Fragos (2020). Mandel & Fragos (2020) reweighted the publicly released χ_{eff} posterior samples from the **EOB** analysis (Abbott et al. 2020) in order to apply a prior that assumes a non-spinning primary and a secondary that has a spin aligned with the orbital angular momentum. This approach assumes that there is a single measurable spin degree of freedom from GW190412 which is identified with χ_{eff} , and that there is no information about spin precession. We instead reanalyze the data under the desired prior, thus imposing no such restrictions about how spins are measured. Our analysis results in similar constraints on the secondary spin (Model D with **EOB**) as Mandel & Fragos (2020), but a different estimate of the mass ratio; we find $0.34 \leq q \leq 0.47$ at the 90% level, compared to $0.27 \leq q \leq 0.36$. This difference could be attributed to the assumptions of Mandel & Fragos (2020) about spin measurability. For example, we find that the data contain small (but non-negligible) information about spin precession. Additionally, the leading-order spin term in the GW phase is not identical to χ_{eff} (Poisson & Will 1995), with the difference between the two being more prominent for unequal mass systems such as GW190412. This suggests that the relation between χ_2 and q cannot be fully explored when considering only χ_{eff} . Both Mandel & Fragos (2020) and this study conclude that

the assumption of a non spinning primary requires a highly spinning secondary, although we find that the corresponding Bayes factors disfavor this scenario.

Regardless of our prior assumptions, we find the positive effective spin and unequal masses of GW190412 to be robust conclusions. However, we do see a shift in the posterior distributions across our prior assumptions. With only the secondary BH spinning, we recover higher values for q and lower values for χ_{eff} and χ_p . The component spins are affected more dramatically; forcing a non-spinning primary causes the secondary’s spin magnitude posterior to significantly increase and rail against the physical boundary at $\chi_2 = 1$.

The sensitivity of parameter-estimation results to the choice of prior highlights the importance of choosing an appropriate prior when interpreting observations. In many situations it is desirable not to use a strong prior from a given model. We will never find a spinning primary BH if we always restrict primary spins to be zero. Astrophysical models are uncertain, and need to be constrained by observations. To this end, we can construct prior distributions using a population of observations. Performing hierarchical inference enables inference of both individual event’s properties and those of the population (Mandel 2010; Abbott et al. 2019b; Galadage et al. 2019), in effect using the set of observations to construct an empirical prior. These inferences may use a branching fraction to consider models different formation channels (Vitale et al. 2017b; Stevenson et al. 2017; Zevin et al. 2017; Talbot & Thrane 2017) or use a phenomenological model to describe the underlying population (Roulet & Zaldarriaga 2018; Fishbach et al. 2018; Wysocki et al. 2019; Fishbach et al. 2020); they may even encode prior odds for different channels (Kimball et al. 2020). Using wide, uninformative priors, as done by the LVC, enables parameter-estimation results to be reweighted by different priors, as required for a hierarchical population analysis (Thrane & Talbot 2019; Mandel et al. 2019).

Both the moderately-spinning primary and highly-spinning secondary interpretations for GW190412 provide unprecedented constraints on astrophysical formation scenarios. If GW190412 is the product of isolated binary evolution, our results indicate that the paradigm of negligible natal spin for the first-born BH in BBH merger progenitors may need to be revised (Kushnir et al. 2016; Zaldarriaga et al. 2018; Fuller & Ma 2019). Recent work has shown that if post-main sequence angular momentum transport is not too strong, the first-born BH in BBH progenitors can be highly spinning from due either a Case-A (main sequence) mass transfer episode or post-main sequence tidal spin-up (Qin et al. 2019).

However, it is unclear if these systems will become BBHs with tight enough orbits to merge within a Hubble time. Alternatively, GW190412 could be of dynamical origin, with the primary BH being the product of one (or more) BBH mergers. The canonical dynamical scenario—formation in a classical globular cluster (Benacquista & Downing 2013)—also struggles to match the parameters of GW190412. To be retained in a globular cluster, the natal spins of first-generation BHs need to be small (e.g., Rodriguez et al. 2019). In this case, the merger product of two BHs will form a second-generation BH with a dimensionless spin of $\chi \approx 0.67$: above the measurement of χ_1 in GW190412, which is $0.09 \leq \chi_1 \leq 0.60$ with the Phenom approximant and $0.31 \leq \chi_1 \leq 0.58$ with the EOB approximant. The second-generation globular cluster scenario for GW190412’s primary BH is also highly disfavored from phenomenological models of hierarchical mergers, which find an odds ratio of $\gtrsim 1000:1$ in favor of a GW190412 being a merger of two first-generation BHs rather than the merger of a first- and second-generation BH in a globular cluster (Kimball et al. 2020). Though globular clusters typically cannot retain higher than second-generation merger products due to the relativistic recoil kicks at merger, nuclear clusters (Gerosa et al. 2020), AGN discs (Tagawa et al. 2020), and high metallicity super star clusters (Rodriguez et al. 2020) have all been proposed for the formation of GW190412 analogs via hierarchically merging BHs. Other more exotic channels have also been proposed for forming GW190412, such as GW190412 resulting from a $3 + 1$ hierarchical quadruple stellar system (Hamers & Safarzadeh 2020), though BBH merger rates from such channels are highly uncertain. Explaining GW observations requires astrophysical models which can produce systems with both parameters and event rates that are consistent with the measured values.

While the particular formation scenario for GW190412 is to be determined, the correct interpretation of GW190412’s component spins (and those of future GW observations) is paramount for constraining viable formation mechanisms. As the GW detector network continues its observational campaign (Abbott et al. 2018), additional observations of asymmetric and spinning systems (or lack thereof) will further inform the astrophysical channels that lead to the formation of merging BBHs.

Posterior samples for the parameter estimation of GW190412 using our suite of prior configurations, as well as model evidences, are available on Zenodo (Zevin et al. 2020).

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